Cross section measurements for $(n,2n),(n,\alpha)$, and (n,p) reactions on rhenium isotopes around 14 MeV neutrons and their theoretical calculations of excitation functions*

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Abstract: Cross-section data of the 185 Re(n,2n) 184 mRe, 185 Re(n,2n) 184 gRe, 185 Re(n,α) $^{182m1+m2+g}$ Ta, 187 Re(n,α) 1

Keywords: cross sections of (n,2n), (n,α) and (n,p) reactions, rhenium isotopes, activation technique, off-line γ -ray spectrometry, theoretical calculations

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I. INTRODUCTION

Accurate and reliable cross sections of nuclear reactions around the neutron energy of 14 MeV on rhenium isotopes are crucial for the design of fusion reactors because rhenium is a critical potential structural component of such a type of reactor [1, 2]. To date, few laboratories (from two to nine) have reported experimental cross-section data of (n,2n), (n,α) , and (n,p) reactions on rhenium isotopes around the neutron energy of 14 MeV. These data can be found in the experimental nuclear reaction data (EXFOR) library [3]. There are large discrepancies in these experimental data that may be due to the different equipment and data processing methods used, interferences of y rays with the same or close energies, neglect of the effect of the excited state on the ground state that should have been deducted, etc. For example, experimental cross-section data of the 185 Re $(n,2n)^{184m}$ Re reaction around the neutron energy of 14 MeV were obtained by nine laboratories [2, 4-11]. There are remarkable dif-

ferences in those data. For instance, the maximum difference between them is given by a factor greater than 3. For the 185 Re $(n,2n)^{184g}$ Re reaction, experimental cross-section data around the neutron energy of 14 MeV were obtained by eight laboratories [2, 4-8, 11, 12]. These data also differ significantly, with a maximum difference greater than 25%. For the 185 Re $(n,\alpha)^{182\text{m1}+\text{m2}+\text{g}}$ Ta reaction, experimental cross-section data were only obtained by two laboratories [4, 13]. These data were suspicious (because of the interference of γ rays with close energies in its cross section measurement) although there are no evident differences between them. For the $^{187}\mathrm{Re}(n,2n)$ the $^{186\mathrm{g},(m)}\mathrm{Re}$ reaction ($^{86\mathrm{g},(m)}\mathrm{Re}$ represents reaction product $^{186\mathrm{g}}\mathrm{Re}$ with isomeric state ^{186m}Re transition contribution), experimental cross-section data around the neutron energy of 14 MeV were obtained by seven laboratories [4, 10-15]. There are notable differences in those data as well, with maximum difference greater than 40%. For the 187 Re $(n,\alpha)^{184}$ Ta reaction, experimental cross-section data around the neutron energy of 14 MeV were obtained by five laboratories [4,

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8, 12, 13, 16], with large differences and a maximum difference given by a factor greater than 2. For the 187 Re $(n,p)^{187}$ W reaction, experimental cross-section data around the neutron energy of 14 MeV were obtained by five laboratories [4, 8, 13, 16, 17], with large differences and a maximum difference greater than 35%. Furthermore, there are differences in the cross-section evaluation values of rhenium isotopes in the neutron energy range from the threshold to 20 MeV given by several maior libraries of the International Atomic Energy Agency (IAEA) [18]. Therefore, additional measurements are required concerning the cross sections of the six reactions mentioned above on rhenium isotopes around the neutron energy of 14 MeV. Besides, their excitation curves must be obtained by using the nuclear theoretical model program system Talys-1.9 [19]. In the present study, the cross sections of the ${}^{185}\text{Re}(n,2n){}^{184\text{m}}\text{Re}, {}^{185}\text{Re}(n,2n){}^{184\text{g}}\text{Re}, {}^{185}\text{Re}(n,2n){}^{184\text{g}}\text{Re}, {}^{185}\text{Re}(n,2n){}^{184\text{g}}\text{Re}, {}^{187}\text{Re}(n,2n){}^{186\text{g},(m)}\text{Re}, {}^{187}\text{Re}(n,2n){}^{184\text{g}}\text{Re}, {}^{$ and $^{187}\text{Re}(n,p)^{187}\text{W}$ reactions were measured around the neutron energies of 13.5-14.8 MeV by using the activation technique and using off-line γ -ray spectrometry. The excitation functions of the above reactions in neutron energies ranging from the threshold to 20 MeV were calculated by adopting the nuclear theoretical model program system Talys-1.9 with the relevant parameters properly adjusted. The measured results were analyzed and compared with previous experiments conducted by other researchers and with the evaluated data of BROND-3.1 (Russia, 2016), ENDF/B-VIII.0 (USA, 2018), and JEFF-3.3 (Europe, 2017) [18], as well as with theoretical values based on Talys-1.9.

II. EXPERIMENTAL DETAILS

Natural rhenium foils of 3.04-3.12 mm thickness and 99.99% purity were made into round disks with a diameter of 20 mm. Natural niobium foils (1 mm in thickness, 99.99% purity) of the same diameter as the rhenium sample were then fixed at the front and back of each rhenium foil, which was wrapped in a cadmium foil of 1 mm thickness and 99.95% purity to reduce the influence of the 185 Re(n, γ) $^{186g,(m)}$ Re reaction induced by low-energy neutrons on the 187 Re(n,2n) $^{186g,(m)}$ Re reaction.

The samples were radiated at the K-400 Neutron Generator at China Academy of Engineering Physics (CAEP) and lasted for 6.5-10 h. Neutrons in the 14 MeV region with a yield ranging from 4×10^{10} n/s to 5×10^{10} n/s were generated from the $T(d,n)^4$ He reaction under a deuteron beam energy of 255 keV and a beam current of 350 μ A. The solid tritium—titanium (T-Ti) target applied to the generator was approximately 2.19 mg/cm² thick. In the process of the radiation of the samples, the neutron flux was detected by an Au-Si surface barrier detector used in 135° accompanying a particle tube to correct small variations of neutron flux. The samples were put at 0°-135°

angles relative to the direction of the deuteron beam and with distances from the center of the T-Ti target of approximately 40-50 mm. The neutron energies in the measurements were determined in advance from the cross section ratios of 90 Zr(n,2n) ${}^{89m+g}$ Zr to 93 Nb(n,2n) 92m Nb reactions [20].

The γ-ray activities of ^{184m}Re, ^{184g}Re, ^{182m1+m2+g}Ta, ^{186g}Re, ¹⁸⁴Ta, ¹⁸⁷W, and ^{92m}Nb were determined by a well-calibrated GEM-60P coaxial high-purity germanium OR-TEC detector made in USA (its crystal diameter is 70.1 mm and crystal length is 72.3 mm) with a relative efficiency of 68% and an energy resolution of 1.69 keV at 1332 keV. The efficiency of the detector was pre-calibrated using various standard γ sources. Activities of decay γ-rays from the product radionuclides were recorded 90 mm away from the detector's surface. The cooling time of the irradiated sample was approximately 18 min-136.32 h, and the measuring time was approximately 23 min-20.78 h. Figures 1 and 2 show part of the γ-ray spectrum obtained from the rhenium samples at approxim-

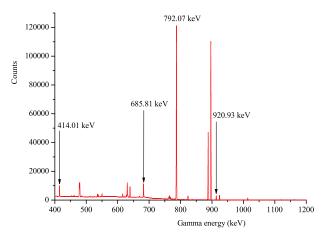


Fig. 1. (color online) Part of the γ -ray spectrum of rhenium obtained after 6.78 h of cooling following the end of irradiation; the measurement duration was approximately 3.82 h.

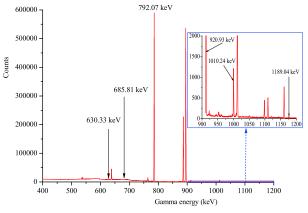


Fig. 2. (color online) Part of the γ -ray spectrum of rhenium obtained after 125.45 h of cooling following the end of irradiation; the measurement duration was approximately 20.78 h.

ately 6.78 and 125.45 h after the end of irradiation, respectively.

III. DATA ANALYSIS

The decay characteristics of the radioactive product nuclides and the natural abundance of the target isotopes under investigation are summarized in Table 1 [21]. The natural abundance of ⁹³Nb was taken from Ref. [22].

The measured cross-section values were calculated by the following formula [23, 24]:

$$\sigma_{x} = \frac{[S \varepsilon I_{\gamma} \eta KMD]_{0}}{[S \varepsilon I_{\gamma} \eta KMD]_{x}} \frac{[\lambda AFC]_{x}}{[\lambda AFC]_{0}} \sigma_{0}. \tag{1}$$

The subscript 0 represents the term corresponding to the monitor reaction, whereas the subscript x corresponds to the measured reaction; ε is the full-energy peak (FEP) efficiency of the measured characteristic γ -ray, I_{γ} is the γ -ray intensity, η is the abundance of the target nuclide, M is the mass of sample, $D = e^{-\lambda t_1} - e^{-\lambda t_2}$ is the counting collection factor, t_1 and t_2 are the time intervals from the end of the irradiation to the start and the end of counting, respectively, A is the atomic weight, C is the measured FEP area, λ is the decay constant, and F is the total correction factor of the activity:

$$F = F_s \times F_c \times F_g, \tag{2}$$

where F_s , F_c , and F_g are correction factors for the self-absorption of the sample at a given γ -energy, the coincidence sum effect of cascade γ -rays of the investigated nuclide, and the counting geometry, respectively. F_c was calculated by the method reported in Ref. [25]. F_s and F_g were calculated by the following equations:

$$F_s = \frac{\mu_m d_m}{1 - e^{-\mu_m d_m}},\tag{3}$$

$$F_g = \frac{(h + d/2)^2}{h^2},\tag{4}$$

where μ_m (in cm²/g) is the mass attenuation coefficient at each gamma energy, d_m (in g/cm²) is the area density, d (in mm) is the thickness of the sample, and h (in mm) is the distance from the surface of the sample to the effective detection cross section of the crystal in the HPGe detector.

K is the neutron fluctuation factor:

$$K = \left[\sum_{i}^{L} \Phi_{i} \left(1 - e^{-\lambda \Delta t_{i}} \right) e^{-\lambda T_{i}} \right] / \Phi S, \tag{5}$$

where L is the number of time intervals into which the irradiation time is divided, Δt_i is the duration of the *i*th time interval, T_i is the time interval from the end of the *i*th interval to the end of irradiation, Φ_i is the neutron flux averaged over the sample during Δt_i , Φ is the neutron flux averaged over the sample during the total irradiation time T, and $S = 1 - e^{-\lambda T}$ is the growth factor of the product nuclide.

Cross sections of the 185 Re(n,2n) 184m Re, 185 Re(n,2n) 184g Re, 185 Re (n,α) ${}^{182m1+m2+g}$ Ta, 187 Re(n,2n) ${}^{186g,(m)}$ Re, 187 Re (n,α) 184 Ta, and 187 Re(n,p) 187 W reactions were obtained. The cross-section data of the monitor reaction 93 Nb $(n,2n)^{92m}$ Nb were 457.9± 6.8, 459.8 ± 6.8, 459.8 ± 6.8, and 459.7 ± 5.0 mb at neutron energies of 13.5, 14.1, 14.4, and 14.8 MeV, respectively [26]. The measured cross-section data are presented in Table 2 and charted in Figs. 3-8. Previously obtained experimental cross sections of these nuclear reactions around the neutron energy of 14 MeV are also summarized in Table 2 and charted in Figs. 3-8 for comparison. Likewise, the evaluation cross-section curves of these reactions from BROND-3.1, ENDF/B-VIII.0, JEFF-3.3, and the theoretical calculation curves in the neutron energy range from the threshold to 20 MeV obtained by the computer code system Talys-1.9 are also charted in Figs. 3-8 for comparison.

To obtain accurate cross-section values of the ${\rm Re}(n,2n)^{\rm 184g}{\rm Re}$ reaction, the effect of the decay of the excited state $^{\rm 184m}{\rm Re}$ on the ground state $^{\rm 184g}{\rm Re}$ and the

 Table 1. Reactions and associated decay data of activation products.

Reaction	Abundance of target isotope (%)	Activation products	$T_{1/2}$	E_{γ}/keV	Ι _γ (%)
185 Re($n,2n$)	37.40	^{184m} Re	169 d	920.93	8.2
185 Re $(n,2n)$	37.40	^{184g} Re	35.4 d	792.07	37.7
				1010.24	0.092
$^{185}\mathrm{Re}(n,\alpha)$	37.40	$^{182m1+m2+g}Ta$	114.74 d	1189.04	16.49
187 Re $(n,2n)$	62.60	^{186g} Re	3.7183 d	630.33	0.0294
$^{187}\mathrm{Re}(n,\alpha)$	62.60	¹⁸⁴ Ta	8.7 h	414.01	72
187 Re (n,p)	62.60	$^{187}{ m W}$	24.0 h	685.81	33.2
93 Nb(<i>n</i> ,2 <i>n</i>)	100	^{92m} Nb	10.15 d	934.44	99.15

Table 2. Summary of the cross sections of Rhenium isotopes around the neutron energy of 14 MeV.

Reaction -	This study		Literature Values		
Reaction	E_n/MeV	σ /mb	E_n/MeV	σ/mb	Reference
185 Re $(n,2n)^{184m}$ Re	13.5±0.2	397±29	14.7	390±18	[2]
	14.1 ± 0.2	409±30	14.74	357±33	[4]
	14.4±0.2	414±30	13.4	403±16	[5]
	14.8±0.2	415±30	13.7	400±16	[5]
			14.1	395±16	[5]
			14.45	389±16	[5]
			14.65	392±16	[5]
			14.8	399±16	[5]
			14.3	380±50	[6]
			14.7	390±70	[6]
			14.44	442±21	[7]
			14.47	428±21	[7]
			13.33	312±26	[8]
			13.57	325±26	[8]
			13.75	332±27	[8]
			13.98	321±26	[8]
			14.22	341±28	[8]
			14.43	335±28	[8]
			14.67	359±30	[8]
			14.93	347±29	[8]
			13.54	322±48	[9]
			13.73	345±52	[9]
			14.67	301±45	[9]
			14.81	416±62	[9]
			14.8	260±100	[10]
			14.1	1120±400	[11]
185 Re $(n,2n)^{184g}$ Re	13.5±0.2	$1800\pm82^{a} (1753\pm155^{b})$	14.7	1817±85	[2]
	14.1±0.2	1787±83 ^a (1727±154 ^b)	14.74	1488±88	[4]
	14.4±0.2	1752±79 ^a (1724±152 ^b)	13.4	1850±74	[5]
	14.8±0.2	1699±77 ^a (1677±149 ^b)	13.7	1844±74	[5
			14.1	1838±74	[5]
			14.45	1832±73	[5]
			14.65	1778±71	[5]
			14.8	1776±71	[5]
			14.3	1670±95	[6]
			14.7	1780±120	[6]
			14.44	1886±75	[7]
			14.47	1886±71	[7]
			13.33	1399±73	[8]
			13.56	1436±75	[8]

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Table 2-continued from previous page

Danction	This work		Literature Values		
Reaction	E_n/MeV	σ/mb	E_n/MeV	σ/mb	Reference
			13.75	1416±74	[8]
			13.98	1389±72	[8]
			14.22	1414±73	[8]
			14.43	1370±72	[8]
			14.66	1385±72	[8]
			14.93	1420±74	[8]
			14.1	1910±600	[11]
			14.7	1900±100	[12]
			15.3	2170±100	[12]
			16.1	1670±90	[12]
			16.8	1260±70	[12]
			17.3	1240±60	[12]
			18.3	900±50	[12]
185 Re $(n,\alpha)^{182m1+m2}$ Ta	13.5±0.2	0.56±0.05	14.48	1.40±0.50	[4]
	14.1±0.2	0.61±0.06	14.7	1.4±0.2	[13]
	14.4±0.2	0.65±0.07			
	14.8±0.2	0.75±0.06			
187 Re $(n,2n)^{186g,(m)}$ Re	13.5±0.2	1734±75	14.85	1680±99	[4]
	14.1±0.2	1661±75	14.8	1490±160	[10]
	14.4±0.2	1614±71	14.1	1440±410	[11]
	14.8±0.2	1560±68	14.7	1140±60	[12]
			15.3	1390±70	[12]
			16.1	950±50	[12]
			16.8	690±40	[12]
			17.3	820±40	[12]
			18.3	710±40	[12]
			13.5	1857±90	[13]
			14.2	1863±90	[13]
			14.7	1881±93	[13]
			14.1	1966.9±50.3	[14]
			14.6	1967.1±44.7	[14]
			14.8	1952.0±50.7	[14]
			15.0	1903.0±50.3	[14]
			14.8	1675±168	[15]
187 Re $(n,\alpha)^{184}$ Ta	13.5±0.2	0.34±0.02	13.67	0.316±0.05	[4]
	14.1±0.2	0.42 ± 0.02	14.06	0.416 ± 0.05	[4]
	14.4±0.2	0.54 ± 0.03	14.46	0.561±0.05	[4]
	14.8±0.2	0.63±0.03	14.84	0.609 ± 0.05	[4]
			14.47	0.445±0.064	[8]
			14.70	0.558±0.097	[8]

Continued on next page

Table 2-continued from previous page

.	This	work		Literature Values	
Reaction	E_n/MeV	σ/mb	E_n/MeV	σ/mb	Reference
			14.96	0.63±0.11	[8]
			14.8	0.53±0.03	[12]
			15.3	0.89 ± 0.05	[12]
			16.1	1.07±0.06	[12]
			16.8	1.06±0.06	[12]
			17.3	0.99 ± 0.05	[12]
			18.3	1.52±0.08	[12]
			14.7	1.1±0.1	[13]
			14.5	0.94 ± 0.14	[16]
187 Re $(n,p)^{187}$ W	13.5±0.2	2.05±0.09	13.67	2.75±0.30	[4]
	14.1 ± 0.2	2.50±0.11	14.06	3.18±0.30	[4]
	14.4 ± 0.2	2.98±0.13	14.46	3.45±0.35	[4]
	14.8±0.2	3.27±0.14	14.84	3.94±0.29	[4]
			13.34	4.51±0.51	[8]
			13.57	4.00±0.71	[8]
			13.76	4.40±0.76	[8]
			13.99	5.10±0.71	[8]
			14.23	4.88±0.79	[8]
			14.44	5.33±0.41	[8]
			14.68	5.59±0.65	[8]
			14.95	6.60±0.75	[8]
			13.5	3±0.3	[13]
			14.2	4±0.4	[13]
			14.7	4.5±0.5	[13]
			14.5	3.93±0.39	[16]
			14.7	4.7±0.4	[17]

^aThese are the results based on the 792.07 keV γ -ray of 184g Re. b These are the results based on the 1010.24 keV γ -ray of 184g Re.

contribution of the 792.07 keV γ -ray (intensity 3.69%) of ^{184m}Re on the 792.07 keV γ -ray (intensity 37.7%) of ^{184g}Re must be deducted. However, the former was not mentioned in Refs. [2, 7] and the latter was not mentioned in Refs. [4, 5]. Our experimental cross-section values of the ¹⁸⁵Re(n,2n) ^{184g}Re reaction were deduced from Equation (1) and 792.07 keV and 1010.24 keV γ -rays of ^{184g}Re, which are listed in Table 2 (marked with superscripts a and b) and charted in Fig. 4. In the process of calculation, the effect of the excited state ^{184m}Re on the ground state ^{184g}Re was deducted by using the method described in Ref. [27]. Simultaneously, the contribution of the 792.07 keV γ -ray of ^{184g}Re was also subtracted by using the method described in Ref. [28].

The theoretical calculations of excitation functions of the $^{185}\mathrm{Re}(n,2n)^{184\mathrm{m}}\mathrm{Re}$, $^{185}\mathrm{Re}(n,2n)^{184\mathrm{g}}\mathrm{Re}$, $^{185}\mathrm{Re}(n,2n)^{184\mathrm{g}}\mathrm{Re}$, $^{185}\mathrm{Re}(n,\alpha)^{184}\mathrm{Ta}$, and $^{187}\mathrm{Re}(n,p)^{187}\mathrm{W}$ reactions were performed by using the nuclear theoretical model program system Talys-1.9. Different parameters in Talys-1.9 were adjusted according to our measured data and previous experiments conducted by other researchers for the different nuclear reactions mentioned above. For the $^{185}\mathrm{Re}(n,2n)^{184\mathrm{m}}\mathrm{Re}$ reaction, the level density parameter at the neutron separation energy was adjusted. For the $^{185}\mathrm{Re}(n,2n)^{184\mathrm{g}}\mathrm{Re}$ reaction, the optical model potential (OMP) parameter $^{185}\mathrm{Re}(n,2n)^{182\mathrm{m}1+\mathrm{m}2+\mathrm{g}}\mathrm{Ta}$, $^{187}\mathrm{Re}(n,2n)^{184}\mathrm{Ta}$, and $^{187}\mathrm{Re}(n,p)^{187}\mathrm{W}$ reactions, the OMP parameter r_V was adjusted. For the

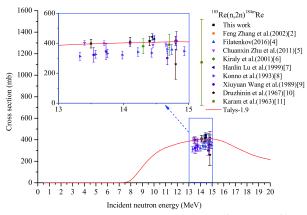


Fig. 3. (color online) Cross section of the 185 Re $(n,2n)^{184m}$ Re reaction.

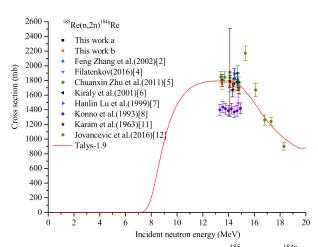


Fig. 4. (color online) Cross section of the 185 Re $(n,2n)^{184g}$ Re reaction.

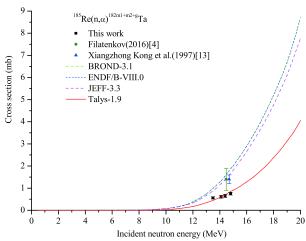


Fig. 5. (color online) Cross section of the 185 Re $(n,\alpha)^{182m1+m2+g}$ Ta reaction.

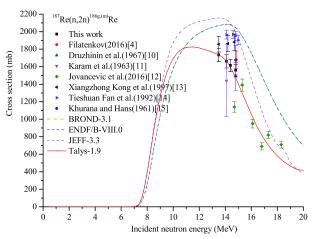


Fig. 6. (color online) Cross section of the 187 Re $(n,2n)^{186g,(m)}$ Re reaction.

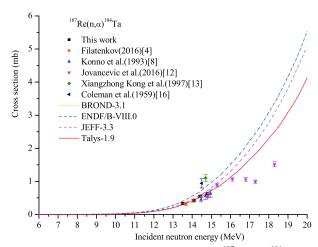


Fig. 7. (color online) Cross section of the $^{187}\text{Re}(n,\alpha)^{184}\text{Ta}$ reaction.

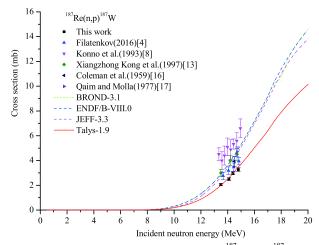


Fig. 8. (color online) Cross section of the $^{187}\text{Re}(n,p)^{187}\text{W}$ reaction.

 $^{187}\mathrm{Re}(n,2n)^{^{186\mathrm{g,(m)}}}\mathrm{Re}$ reaction, the OMP parameter r_V and the model for level densities were adjusted.

IV. RESULTS AND DISCUSSION

The principal sources of uncertainty and their estimated values in the present study are given in Table 3.

For the 185 Re(n,2n) Re reaction, note from Table 2 and Fig. 3 that the theoretical excitation curve based on Talys-1.9 code increases with increasing neutron energy in the incident neutron energy range from the threshold to 15 MeV but decreases with increasing neutron energy in the incident neutron energy range from 15 to 20 MeV. Our experimental cross-section values are consistent, within experimental error, with those of the fitting line of the results of Chuanxin Zhu et al. [5] and the theoretical excitation curve based on Talys-1.9 code at the corresponding energies. The results reported by Kiraly et al. [6] and the value reported by Xiuyuan Wang et al. [9] at the neutron energy of 14.81 MeV, within experimental error, are consistent with those of the fitting line of the results of Chuanxin Zhu et al. [5] and our experimental values at the corresponding energies, and with those of the theoretical excitation curve based on Talys-1.9 code at the corresponding energies. In contrast, the results reported by Konno et al. [8], the values of Xiuyuan Wang et al. [9] at the neutron energies of 13.54, 13.73, and 14.67 MeV, and the value reported by Druzhinin et al. [10] are lower, whereas the cross-section values reported by Hanlin Lu et al. [7] and the results reported by Karam et al. [11] are higher.

For the 185 Re(n,2n) 184g Re reaction, Table 2 and Fig. 4 show that the results obtained in this study by using the two aforementioned γ -rays of 184g Re are highly consistent, within experimental error; the theoretical excitation curve based on Talys-1.9 code increases with increasing neutron energy in the incident neutron energy range from the

Table 3. Principal sources of uncertainty and their estimated values in the cross section measurements.

Source of uncertainty	Uncertainty (%)		
γ-ray detection efficiency	2		
Standard cross section	1.1-1.5		
Counting statistics	0.2-9.7		
Weight of samples	0.1		
Sample geometry	1.0		
Self-absorption of γ ray	1.0		
Neutron fluctuation	1.0		
Relative γ -ray intensity	0.3-7.6		
Half-life	0.02-4.7		
Isotopic abundance	0.03-0.05		
Total uncertainty	4.03-10.5		

threshold to 13 MeV but decreases with increasing neutron energy in the incident neutron energy range from 13 to 20 MeV, thereby matching most of the experimental data well. Our experimental cross-section values are consistent, within experimental error, with those of the fitting line of the results of Chuanxin Zhu *et al.* [5] and the theoretical excitation curve based on Talys-1.9 code at the corresponding energies. In contrast, the results reported by Hanlin Lu *et al.* [7] and those of Jovancevic *et al.* [12] at the neutron energies of 14.7, 15.3, and 16.1 MeV are higher, whereas the results reported by Filatenkov [4] and the cross-section values reported by Konno *et al.* [8] are lower.

For the 185 Re (n,α) $^{182\text{m1+m2+g}}$ Ta reaction, as shown in Table 2 and Fig. 5, three evaluated excitation curves of BROND-3.1, ENDF/B-VIII.0, JEFF-3.3 and the theoretical excitation curve based on Talys-1.9 code increase with neutron energy in the incident neutron energy range from the threshold to 20 MeV, but there are differences between them, except for the curves of BROND-3.1 and ENDF/B-VIII.0 (which are the same). The three evaluated excitation curves pass through the experimental point of Filatenkov [4], and the curve of JEFF-3.3 passes through the experimental points of Filatenkov [4] and Xiangzhong Kong et al. [13]. The theoretical excitation curve passes through our four experimental points. The experimental values reported by Filatenkov [4] and Xiangzhong Kong et al. [13] are higher than those of the fitting line of our experimental values and the theoretical excitation curve obtained by Talvs-1.9 at the corresponding energies. A possible reason is that the interactions of γ -rays with the same or close energies should be avoided or deducted but not avoided or deducted, such as the cross-section value of the $^{185}\mathrm{Re}(n,\alpha)^{^{182\mathrm{ml+m2+g}}}\mathrm{Ta}$ reaction was deducted by the 1189.0 and 1221.4 keV γ-rays of ^{182g}Ta (half life 114.74 d) in Ref.[4], in the process of calculation. Likewise, the contribution of the 1121.29 keV (intensity 0.0202%) γ -ray of 184g Re (from the Re $(n,2n)^{184g}$ Re reaction) on the full-energy peak (FEP) count of the 1221.4 keV γ -ray of Ta should be deducted, which was not mentioned in Ref.[4]. For the 187 Re(n,2n) $^{186g,(m)}$ Re reaction, shown in Table 2

For the Re(*n*,2*n*) residual Re reaction, shown in Table 2 and Fig. 6, the trends of the evaluated excitation curves of BROND-3.1, ENDF/B-VIII.0, JEFF-3.3, and the theoretical excitation curve based on Talys-1.9 code decrease with increasing neutron energy around 14 MeV. However, the extent of their reduction is different, except for the curves of BROND-3.1 and ENDF/B-VIII.0 (which are the same). Our experimental cross-section values are consistent, within experimental error, with those of Druzhinin *et al.* [10], Karam *et al.* [11], and Khurana and Hans [15] at the same energies, and with those of the theoretical excitation curve based on Talys-1.9 code at the corresponding energies besides the neutron energy of 14.8 MeV. The value reported by Filatenkov [4] is con-

sistent, within experimental error, with that of the fitting line of our experimental cross-section values at the corresponding energy. Note that the cross-section values of Xiangzhong Kong et al. [13] were deduced from Equation (1) and the 137.14 keV γ -ray (intensity 8.5%) of ^{186g}Re (half life 3.777 d), whose latest value is 137.2 keV (intensity 9.47%) of ^{186g}Re (half life 3.7183 d). The cross section values corrected with the latest parameter were 1693, 1698, and 1714 mb at the neutron energies of 13.5, 14.2, and 14.7 MeV, respectively, which are consistent with those of the fitting line of our experimental crosssection values, within experimental error. The results reported by Fan Tieshuan et al. [14] are considerably higher than the other experimental values (including ours) and those of the theoretical excitation curve at the corresponding energies. Compared with the three evaluation curves, the theoretical curve based on Talys-1.9 code is in better agreement with the experimental results.

For the 187 Re $(n,\alpha)^{184}$ Ta reaction, shown in Table 2 and Fig. 7, the trends of the three evaluated excitation curves of BROND-3.1, ENDF/B-VIII.0, JEFF-3.3, and the theoretical excitation curve based on Talys-1.9 code increase with increasing neutron energy in the incident neutron energy range from the threshold to 20 MeV, but there are differences between them, except for the curves of BROND-3.1 and ENDF/B-VIII.0 (which are the same). Our experimental cross-section values are consistent, within experimental error, with those of the fitting lines of the results reported by Filatenkov [4] and the corrected cross-section values reported by Konno et al. [8], which were corrected with the latest intensity, 72%, of the 414.01 keV γ -ray of ¹⁸⁴Ta instead of 74%, and with those of the theoretical excitation curve based on Talys-1.9 code at the corresponding energies. Comparison with the three evaluation curves shows that the obtained theoretical excitation curve based on Talys-1.9 code can match most of the experimental data. In contrast, the values of Jovancevic et al. [12] at the neutron energies of 16.8, 17.3, and 18.3 MeV are lower, and the values of Xiangzhong Kong et al. [13] and Coleman et al. [16] are high-

For the ¹⁸⁷Re(*n*,*p*)¹⁸⁷W reaction, shown in Table 2 and Fig. 8, the trends of the three evaluated excitation curves of BROND-3.1, ENDF/B-VIII.0, JEFF-3.3, and the theoretical excitation curve based on Talys-1.9 code increase with increasing neutron energy in the incident neutron energy range from the threshold to 20 MeV, but there are differences between them, except for the curves of BROND-3.1 and ENDF/B-VIII.0 (which are the same). Our experimental cross-section values are consistent, within experimental error, with those of the theoretic-

al excitation curve based on Talys-1.9 code at the corresponding energies. These previous experimental values and those of the three evaluated excitation curves at the corresponding energies are all significantly higher than our experimental cross-section values. The reason may be that the cross-section values of the $^{187}\text{Re}(n,p)^{187}\text{W}$ reaction were deduced by the 479.53 keV (intensity 26.6%) γ ray of ^{187}W (half life 24.0 h) in Refs. [4, 8, 13, 17], and in the process of calculation, the contribution of the 478.0 keV (intensity 1.017%) γ ray of ^{188g}Re (from the ^{187}Re (n,γ) ^{188g}Re reaction) on the FEP count of the 479.53 keV γ ray of ^{187}W should be deducted but not deducted; moreover, the data in Ref. [16] were obtained via β-ray counting.

V. CONCLUSION

cross-section data Experimental of the 185 Re $(n,2n)^{^{184m}}$ Re $^{^{185}}$ Re $(n,2n)^{^{184g}}$ Re $^{^{185}}$ Re $(n,\alpha)^{^{182m1+m2+g}}$ Ta, 187 Re $(n,2n)^{^{186g,(m)}}$ Re $^{^{187}}$ Re $(n,\alpha)^{^{184}}$ Ta, and 187 Re $(n,p)^{^{187}}$ W reactions were measured in the neutron energy range of 13.5-14.8 MeV via the activation technique. The excitation functions of the aforementioned six reactions in the neutron energies from the threshold to 20 MeV were calculated by adopting the nuclear theoretical model program system Talys-1.9 with the relevant parameters properly adjusted. The measured cross sections were discussed and compared with previous experiments by other researchers, and with the evaluated data of BROND-3.1, ENDF/B-VIII.0, JEFF-3.3, and the theoretical results based on Talys-1.9 code. In general, our experimental cross-section values are consistent, within experimental error, with those of previous experiments and theoretical excitation curves at the corresponding energies. Comparison with these evaluation curves shows that the theoretical excitation curves based on Talys-1.9 code agree well with the experimental results. The new measured results in the present study would improve the quality of the neutron cross section database and are expected to assist with new evaluations of cross sections on rhenium isotopes in the incident neutron energy range from the threshold to 20 MeV. In addition, the theoretical excitation curves are relevant for the design of fusion reactors and related applications.

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